### **Optics Letters**

## Probing coherence Stokes parameters of three-component light with nanoscatterers: supplement

MENGWEN GUO, 1,2,\* D ANDREAS NORRMAN, 2,3 ARI T. FRIBERG, 2 AND TERO SETÄLÄ<sup>2</sup>

This supplement published with Optica Publishing Group on 11 May 2022 by The Authors under the terms of the Creative Commons Attribution 4.0 License in the format provided by the authors and unedited. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI.

Supplement DOI: https://doi.org/10.6084/m9.figshare.19619865

Parent Article DOI: https://doi.org/10.1364/OL.457963

<sup>&</sup>lt;sup>1</sup>Department of Physics, Hangzhou Dianzi University, Hangzhou 310018, China

<sup>&</sup>lt;sup>2</sup>Institute of Photonics, University of Eastern Finland, P.O. Box 111, FI-80101 Joensuu, Finland

<sup>&</sup>lt;sup>3</sup>Photonics Laboratory, ETH Zurich, CH-8093 Zurich, Switzerland

<sup>\*</sup>Corresponding author: mengwen.guo@uef.fi

# Probing coherence Stokes parameters of three-component light with nanoscatterers: supplemental document

This supplementary document provides detailed derivations of the difference between the distances from the dipoles to the far-zone point as well as Eqs. (16a)–(16i) appearing in the main text

#### **DERIVATION OF DISTANCE DIFFERENCE**

Suppose that the two dipoles are located symmetrically about the origin along the x axis with separation d=2a (see Fig. 1). In the Cartesian coordinate frame we then have  $\mathbf{r}_1=(-a,0,0)$ ,  $\mathbf{r}_2=(a,0,0)$ , and  $\mathbf{r}=(r\sin\theta\cos\varphi,r\sin\theta\sin\varphi,r\cos\theta)$ , where  $r=|\mathbf{r}|$ . Hence,

$$r_{1} = \sqrt{(r\sin\theta\cos\varphi + a)^{2} + (r\sin\theta\sin\varphi)^{2} + (r\cos\theta)^{2}}$$

$$= r\sqrt{1 + \left(\frac{a}{r}\right)^{2} + 2\frac{a}{r}\sin\theta\cos\varphi}.$$
(S1)

Since  $r \gg a$ , we can neglect the second term and use  $(1+x)^{1/2} \approx 1 + x/2$  to obtain

$$r_1 \approx r \left(1 + \frac{a}{r} \sin \theta \cos \varphi\right).$$
 (S2)

In a similar way we find that

$$r_{2} = \sqrt{(r\sin\theta\cos\varphi - a)^{2} + (r\sin\theta\sin\varphi)^{2} + (r\cos\theta)^{2}}$$

$$\approx r(1 - \frac{a}{r}\sin\theta\cos\varphi). \tag{S3}$$

Equations (S2) and (S3) imply at once that

$$k(r_1 - r_2) \approx 2ka\sin\theta\cos\varphi = kd\sin\theta\cos\varphi,$$
 (S4)

which was employed in the discussion above Eq. (14) of the main text.

### **DERIVATION OF EQUATIONS (16a)–(16i)**

The three observation directions  $(\theta, \varphi)$  under consideration are  $(\pi/2, 0)$ ,  $(\pi/2, \pi/2)$ , and  $(0, \pi/2)$ , as illustrated in Fig. 2.

Direction  $(\pi/2,0)$ : From Eqs. (5) and (6) one finds that  $E'_{\theta}(\mathbf{r}_i,\omega)=-E_z(\mathbf{r}_i,\omega)$  and  $E'_{\varphi}(\mathbf{r}_i,\omega)=E_y(\mathbf{r}_i,\omega)$ , which are the field components at the dipole location  $\mathbf{r}_i$  along the  $\hat{\mathbf{u}}_{\theta}$  and  $\hat{\mathbf{u}}_{\varphi}$  directions, respectively. Thus the elements of the dipole-site  $2\times 2$  cross-spectral density matrix constructed from the components transverse to the direction  $(\pi/2,0)$  are related to the Cartesian-system elements as

$$\begin{bmatrix} W_{\theta\theta}^{\prime(12)} & W_{\theta\varphi}^{\prime(12)} \\ W_{\varphi\theta}^{\prime(12)} & W_{\varphi\varphi}^{\prime(12)} \end{bmatrix} = \begin{bmatrix} W_{zz}^{(12)} & -W_{zy}^{(12)} \\ -W_{yz}^{(12)} & W_{yy}^{(12)} \end{bmatrix}, \tag{S5}$$

where the superscript (12) refers to points  $\mathbf{r}_1$  and  $\mathbf{r}_2$ . Using Eqs. (13a), (13c), and (13d) then yields

$$S_0^{\prime(12)}(\frac{\pi}{2},0) = W_{zz}^{(12)} + W_{yy}^{(12)},$$
 (S6)

$$S_2^{\prime(12)}(\frac{\pi}{2},0) = -[W_{yz}^{(12)} + W_{zy}^{(12)}],$$
 (S7)

$$S_3^{\prime(12)}(\frac{\pi}{2},0) = i[W_{yz}^{(12)} - W_{zy}^{(12)}].$$
 (S8)

On further substituting Eqs. (S7) and (S8) into Eqs. (2g) and (2h), respectively, one finds

$$\Lambda_6^{(12)} = -S_2^{\prime(12)}(\frac{\pi}{2}, 0), \tag{S9}$$

$$\Lambda_7^{(12)} = S_3^{\prime(12)}(\frac{\pi}{2}, 0),\tag{S10}$$

which are Eqs. (16g) and (16h) of the main text.  $Direction~(\pi/2,\pi/2)$ : Now  $E'_{\theta}(\mathbf{r}_i,\omega)=-E_z(\mathbf{r}_i,\omega)$  and  $E'_{\varphi}(\mathbf{r}_i,\omega)=-E_x(\mathbf{r}_i,\omega)$ , whereupon

$$\begin{bmatrix} W_{\theta\theta}^{\prime(12)} & W_{\theta\varphi}^{\prime(12)} \\ W_{\varphi\theta}^{\prime(12)} & W_{\varphi\varphi}^{\prime(12)} \end{bmatrix} = \begin{bmatrix} W_{zz}^{(12)} & W_{zx}^{(12)} \\ W_{xz}^{(12)} & W_{xx}^{(12)} \end{bmatrix}.$$
 (S11)

The use of Eqs. (13a), (13c), and (13d) subsequently results in

$$S_0^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}) = W_{zz}^{(12)} + W_{xx}^{(12)},$$
 (S12)

$$S_2^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}) = W_{xz}^{(12)} + W_{zx}^{(12)},$$
 (S13)

$$S_3^{\prime(12)}(\frac{\pi}{2},\frac{\pi}{2}) = -i[W_{xz}^{(12)} - W_{zx}^{(12)}]. \tag{S14}$$

By then inserting Eqs. (S13) and (S14) into Eqs. (2e) and (2f), respectively, one obtains

$$\Lambda_4^{(12)} = S_2^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}),\tag{S15}$$

$$\Lambda_5^{(12)} = -S_3^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}),\tag{S16}$$

which are Eqs. (16e) and (16f) of the body text.  $Direction~(0,\pi/2)$ : In this case  $E'_{\theta}(\mathbf{r}_i,\omega)=E_y(\mathbf{r}_i,\omega)$  and  $E'_{\phi}(\mathbf{r}_i,\omega)=-E_x(\mathbf{r}_i,\omega)$ . Therefore,

$$\begin{bmatrix} W_{\theta\theta}^{\prime(12)} & W_{\theta\varphi}^{\prime(12)} \\ W_{\varphi\theta}^{\prime(12)} & W_{\varphi\varphi}^{\prime(12)} \end{bmatrix} = \begin{bmatrix} W_{yy}^{(12)} & -W_{yx}^{(12)} \\ -W_{xy}^{(12)} & W_{xx}^{(12)} \end{bmatrix}, \tag{S17}$$

and following a similar procedure as above leads to

$$S_0^{\prime(12)}(0,\frac{\pi}{2}) = W_{xx}^{(12)} + W_{yy}^{(12)},$$
 (S18)

$$S_2^{\prime(12)}(0,\frac{\pi}{2}) = -[W_{xy}^{(12)} + W_{yx}^{(12)}],$$
 (S19)

$$S_3^{\prime(12)}(0,\frac{\pi}{2}) = i[W_{xy}^{(12)} - W_{yx}^{(12)}].$$
 (S20)

Comparing Eqs. (S19) and (S20) with Eqs. (2b) and (2c) finally reveals that

$$\Lambda_1^{(12)} = -S_2^{\prime(12)}(0, \frac{\pi}{2}),\tag{S21}$$

$$\Lambda_2^{(12)} = S_3^{\prime(12)}(0, \frac{\pi}{2}),\tag{S22}$$

which are Eqs. (16b) and (16c) of the paper.

Remaining relations: From Eqs. (2a), (2d), (2i) and (S6), (S12), (S18) one straightforwardly obtains

$$\Lambda_0^{(12)} = \frac{1}{2} \left[ S_0^{\prime(12)}(\frac{\pi}{2}, 0) + S_0^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}) + S_0^{\prime(12)}(0, \frac{\pi}{2}) \right], \tag{S23}$$

$$\Lambda_3^{(12)} = S_0^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}) - S_0^{\prime(12)}(\frac{\pi}{2}, 0), \tag{S24}$$

$$\Lambda_8^{(12)} = \frac{1}{\sqrt{3}} \left[ 2S_0^{\prime(12)}(0, \frac{\pi}{2}) - S_0^{\prime(12)}(\frac{\pi}{2}, \frac{\pi}{2}) - S_0^{\prime(12)}(\frac{\pi}{2}, 0) \right], \tag{S25}$$

which are Eqs. (16a), (16d), and (16i).