# Probing coherence Stokes parameters of three-component light with nanoscatterers: supplement 

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## Probing coherence Stokes parameters of three-component light with nanoscatterers: supplemental document

This supplementary document provides detailed derivations of the difference between the distances from the dipoles to the far-zone point as well as Eqs. (16a)-(16i) appearing in the main text.

## DERIVATION OF DISTANCE DIFFERENCE

Suppose that the two dipoles are located symmetrically about the origin along the $x$ axis with separation $d=2 a$ (see Fig. 1). In the Cartesian coordinate frame we then have $\mathbf{r}_{1}=(-a, 0,0)$, $\mathbf{r}_{2}=(a, 0,0)$, and $\mathbf{r}=(r \sin \theta \cos \varphi, r \sin \theta \sin \varphi, r \cos \theta)$, where $r=|\mathbf{r}|$. Hence,

$$
\begin{align*}
r_{1} & =\sqrt{(r \sin \theta \cos \varphi+a)^{2}+(r \sin \theta \sin \varphi)^{2}+(r \cos \theta)^{2}} \\
& =r \sqrt{1+\left(\frac{a}{r}\right)^{2}+2 \frac{a}{r} \sin \theta \cos \varphi .} \tag{S1}
\end{align*}
$$

Since $r \gg a$, we can neglect the second term and use $(1+x)^{1 / 2} \approx 1+x / 2$ to obtain

$$
\begin{equation*}
r_{1} \approx r\left(1+\frac{a}{r} \sin \theta \cos \varphi\right) . \tag{S2}
\end{equation*}
$$

In a similar way we find that

$$
\begin{align*}
r_{2} & =\sqrt{(r \sin \theta \cos \varphi-a)^{2}+(r \sin \theta \sin \varphi)^{2}+(r \cos \theta)^{2}} \\
& \approx r\left(1-\frac{a}{r} \sin \theta \cos \varphi\right) . \tag{S3}
\end{align*}
$$

Equations (S2) and (S3) imply at once that

$$
\begin{equation*}
k\left(r_{1}-r_{2}\right) \approx 2 k a \sin \theta \cos \varphi=k d \sin \theta \cos \varphi \tag{S4}
\end{equation*}
$$

which was employed in the discussion above Eq. (14) of the main text.

## DERIVATION OF EQUATIONS (16a)-(16i)

The three observation directions $(\theta, \varphi)$ under consideration are $(\pi / 2,0),(\pi / 2, \pi / 2)$, and $(0, \pi / 2)$, as illustrated in Fig. 2.
Direction ( $\pi / 2,0)$ : From Eqs. (5) and (6) one finds that $E_{\theta}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=-E_{z}\left(\mathbf{r}_{i}, \omega\right)$ and $E_{\varphi}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=$ $E_{y}\left(\mathbf{r}_{i}, \omega\right)$, which are the field components at the dipole location $\mathbf{r}_{i}$ along the $\hat{\mathbf{u}}_{\theta}$ and $\hat{\mathbf{u}}_{\varphi}$ directions, respectively. Thus the elements of the dipole-site $2 \times 2$ cross-spectral density matrix constructed from the components transverse to the direction $(\pi / 2,0)$ are related to the Cartesian-system elements as

$$
\left[\begin{array}{ll}
W_{\theta \theta}^{\prime(12)} & W_{\theta \varphi}^{\prime(12)}  \tag{S5}\\
W_{\varphi \theta}^{\prime(12)} & W_{\varphi \varphi}^{\prime(12)}
\end{array}\right]=\left[\begin{array}{cc}
W_{z z}^{(12)} & -W_{z y}^{(12)} \\
-W_{y z}^{(12)} & W_{y y}^{(12)}
\end{array}\right],
$$

where the superscript (12) refers to points $\mathbf{r}_{1}$ and $\mathbf{r}_{2}$. Using Eqs. (13a), (13c), and (13d) then yields

$$
\begin{align*}
& S_{0}^{\prime(12)}\left(\frac{\pi}{2}, 0\right)=W_{z z}^{(12)}+W_{y y}^{(12)}  \tag{S6}\\
& S_{2}^{\prime(12)}\left(\frac{\pi}{2}, 0\right)=-\left[W_{y z}^{(12)}+W_{z y}^{(12)}\right]  \tag{S7}\\
& S_{3}^{\prime(12)}\left(\frac{\pi}{2}, 0\right)=\mathrm{i}\left[W_{y z}^{(12)}-W_{z y}^{(12)}\right] \tag{S8}
\end{align*}
$$

On further substituting Eqs. (S7) and (S8) into Eqs. (2g) and (2h), respectively, one finds

$$
\begin{align*}
& \Lambda_{6}^{(12)}=-S_{2}^{\prime(12)}\left(\frac{\pi}{2}, 0\right),  \tag{S9}\\
& \Lambda_{7}^{(12)}=S_{3}^{\prime(12)}\left(\frac{\pi}{2}, 0\right), \tag{S10}
\end{align*}
$$

which are Eqs. $(16 \mathrm{~g})$ and $(16 \mathrm{~h})$ of the main text.
Direction $(\pi / 2, \pi / 2)$ : Now $E_{\theta}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=-E_{z}\left(\mathbf{r}_{i}, \omega\right)$ and $E_{\varphi}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=-E_{x}\left(\mathbf{r}_{i}, \omega\right)$, whereupon

$$
\left[\begin{array}{ll}
W_{\theta \theta}^{\prime(12)} & W_{\theta \varphi}^{\prime(12)}  \tag{S11}\\
W_{\varphi \theta}^{\prime(12)} & W_{\varphi \varphi}^{\prime(12)}
\end{array}\right]=\left[\begin{array}{ll}
W_{z z}^{(12)} & W_{z x}^{(12)} \\
W_{x z}^{(12)} & W_{x x}^{(12)}
\end{array}\right] .
$$

The use of Eqs. (13a), (13c), and (13d) subsequently results in

$$
\begin{align*}
& S_{0}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)=W_{z z}^{(12)}+W_{x x}^{(12)}  \tag{S12}\\
& S_{2}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)=W_{x z}^{(12)}+W_{z x}^{(12)}  \tag{S13}\\
& S_{3}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)=-\mathrm{i}\left[W_{x z}^{(12)}-W_{z x}^{(12)}\right] . \tag{S14}
\end{align*}
$$

By then inserting Eqs. (S13) and (S14) into Eqs. (2e) and (2f), respectively, one obtains

$$
\begin{align*}
& \Lambda_{4}^{(12)}=S_{2}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)  \tag{S15}\\
& \Lambda_{5}^{(12)}=-S_{3}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right), \tag{S16}
\end{align*}
$$

which are Eqs. (16e) and (16f) of the body text.
Direction $(0, \pi / 2)$ : In this case $E_{\theta}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=E_{y}\left(\mathbf{r}_{i}, \omega\right)$ and $E_{\varphi}^{\prime}\left(\mathbf{r}_{i}, \omega\right)=-E_{x}\left(\mathbf{r}_{i}, \omega\right)$. Therefore,

$$
\left[\begin{array}{ll}
W_{\theta \theta}^{\prime(12)} & W_{\theta \varphi}^{\prime(12)}  \tag{S17}\\
W_{\varphi \theta}^{\prime(12)} & W_{\varphi \varphi}^{\prime(12)}
\end{array}\right]=\left[\begin{array}{cc}
W_{y y}^{(12)} & -W_{y x}^{(12)} \\
-W_{x y}^{(12)} & W_{x x}^{(12)}
\end{array}\right],
$$

and following a similar procedure as above leads to

$$
\begin{align*}
& S_{0}^{\prime(12)}\left(0, \frac{\pi}{2}\right)=W_{x x}^{(12)}+W_{y y}^{(12)},  \tag{S18}\\
& S_{2}^{\prime(12)}\left(0, \frac{\pi}{2}\right)=-\left[W_{x y}^{(12)}+W_{y x}^{(12)}\right],  \tag{S19}\\
& S_{3}^{\prime(12)}\left(0, \frac{\pi}{2}\right)=\mathrm{i}\left[W_{x y}^{(12)}-W_{y x}^{(12)}\right] . \tag{S20}
\end{align*}
$$

Comparing Eqs. (S19) and (S20) with Eqs. (2b) and (2c) finally reveals that

$$
\begin{align*}
& \Lambda_{1}^{(12)}=-S_{2}^{\prime(12)}\left(0, \frac{\pi}{2}\right),  \tag{S21}\\
& \Lambda_{2}^{(12)}=S_{3}^{\prime(12)}\left(0, \frac{\pi}{2}\right), \tag{S22}
\end{align*}
$$

which are Eqs. (16b) and (16c) of the paper.
Remaining relations: From Eqs. (2a), (2d), (2i) and (S6), (S12), (S18) one straightforwardly obtains

$$
\begin{align*}
& \Lambda_{0}^{(12)}=\frac{1}{2}\left[S_{0}^{\prime(12)}\left(\frac{\pi}{2}, 0\right)+S_{0}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)+S_{0}^{\prime(12)}\left(0, \frac{\pi}{2}\right)\right],  \tag{S23}\\
& \Lambda_{3}^{(12)}=S_{0}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)-S_{0}^{\prime(12)}\left(\frac{\pi}{2}, 0\right),  \tag{S24}\\
& \Lambda_{8}^{(12)}=\frac{1}{\sqrt{3}}\left[2 S_{0}^{\prime(12)}\left(0, \frac{\pi}{2}\right)-S_{0}^{\prime(12)}\left(\frac{\pi}{2}, \frac{\pi}{2}\right)-S_{0}^{\prime(12)}\left(\frac{\pi}{2}, 0\right)\right], \tag{S25}
\end{align*}
$$

which are Eqs. (16a), (16d), and (16i).

